

# THEORY OF CURRENT AND VOLTAGE RESONANCE IN RLC CIRCUIT AND THEIR MILITARY SIGNIFICANCE

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## Abstract

This article is devoted to current and voltage resonances in RLC oscillatory circuits. Particular attention is paid to the theoretical foundations of resonant processes.

**Keywords:** Resonance, oscillatory circuit, charge resonance, current resonance, voltage resonance, quality factor, differential equation, driving force, damped oscillation, damping coefficient, vector diagram.

## Introduction

**Аннотация:** данная статья посвящена резонансам тока и напряжений в RLC-колебательных цепях. Особое внимание уделено теоретическим основам резонансных процессов.

**Ключевые слова:** резонанс, колебательный контур, резонанс заряда, резонанс тока, резонанс напряжений, добротность, дифференциальное уравнение, вынуждающая сила, затухающее колебание, коэффициент затухания, векторная диаграмма.

If the value of the oscillatory circuit is great in the formation, or more precisely, in the generation of electromagnetic oscillations and waves, then the value of the resonance phenomenon in the oscillatory circuit is infinite in the transmission of electromagnetic waves into space or the capture and use of electromagnetic waves. This is the basis of the principle of radio transmission and reception [1]. Based on this, thanks to this article we will get acquainted with the theory of resonances of current and voltage in RLC oscillatory circuits.

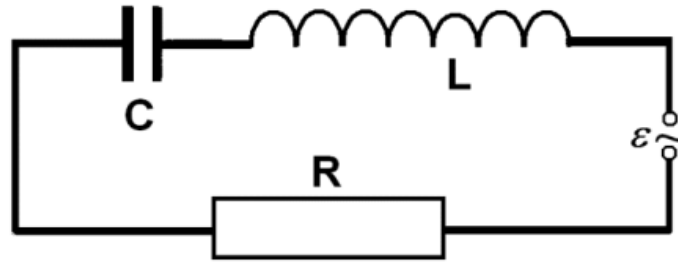


Fig. 1. RLC oscillatory circuit diagram

Based on Figure 1, we will write down Kirchhoff's second law for this circuit to consider the theory of the process occurring in the RLC oscillatory circuit under the influence of variable EMF:  $e = e_0 \cos \omega t$ :

$$R \cdot I + U_C = e - L \frac{dI}{dt} \quad (1)$$

Here  $U_C = q/C$  is the voltage drop across the capacitor,  $I = dq/dt$  is the current in the circuit. Let us write equation (1) as a differential equation:

$$L \frac{d^2q}{dt^2} + R \cdot \frac{dq}{dt} + \frac{q}{C} = e_0 \cos \omega t \quad (2)$$

Dividing all terms of the equation by  $L$ , we obtain a differential equation with a non-homogeneous constant coefficient for the charge  $q$ :

$$\ddot{q} + \frac{R}{L} \cdot \dot{q} + \frac{q}{LC} = \frac{e_0}{L} \cos \omega t \quad (3)$$

From here we obtain a differential equation in which both damping and driving forces participate:

$$\ddot{q} + 2\beta\dot{q} + \omega_0^2 q = f_0 \cos \omega t \quad (4)$$

Where  $\beta = \frac{R}{2L}$ ,  $\omega_0 = \sqrt{\frac{1}{LC}}$ ,  $f_0 = \frac{e_0}{L}$ , As we know, mass is a measure of inertia in mechanical processes, but in electromagnetic processes the measure of inertia is inductance, and the reciprocal value of capacitance plays the role of the coefficient of elasticity in a spring pendulum.

The general solution of such a differential equation is sought as the sum of the general solution of the homogeneous equation  $x_1$  and the particular solution of the inhomogeneous equation  $x_2$ :

$$x = x_{1, \text{general homogeneous}} + x_{2, \text{private heterogeneous}};$$

$$x = x_{1, \text{general homogeneous}} + x_{2, \text{damped oscillation}} = A_1 e^{-\beta t} \sin(\omega_1 t + \alpha_1);$$

Majburlovchi kuch bo‘lmaganda, ya’ni o‘zgaruvcha EYuK bo‘lmagandagi so‘nuvchi tebranish chastotasi ifodasi quyidagicha bo‘ladi:

$$\omega_1 = \sqrt{\omega_0^2 - \beta^2} = \sqrt{\frac{1}{LC} - \left(\frac{R}{2L}\right)^2}$$

The frequency of free oscillations of the system in the absence of a driving force, and in the case of an RLC circuit – in the absence of a variable EMF:

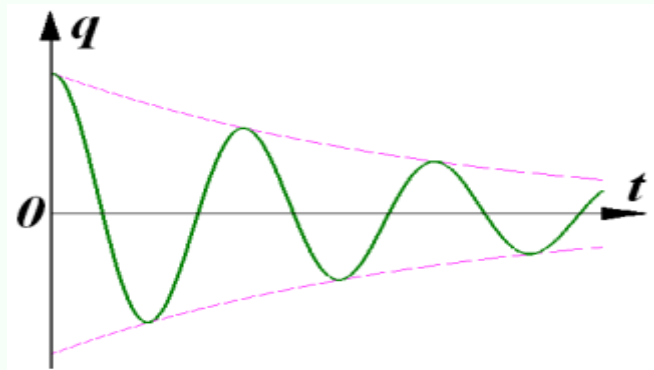


Fig. 2. Graph of free electromagnetic oscillation

The solution  $x_2$  (2, partial non-uniform) defines in mechanics the nature of the motion of a system that performs steady-state forced oscillations, and for an oscillatory circuit  $x_2$  (2, partial non-uniform) describes the law of change of charge on the capacitor plates under the condition of action of EMF in the oscillatory circuit  $e = e_0 \cos \omega t$  [3].

To find  $x_{2, \text{damped oscillation}} = A_2 \cos(\omega t - \varphi_2)$  we use a vector diagram as it was done when studying mechanical vibrations:

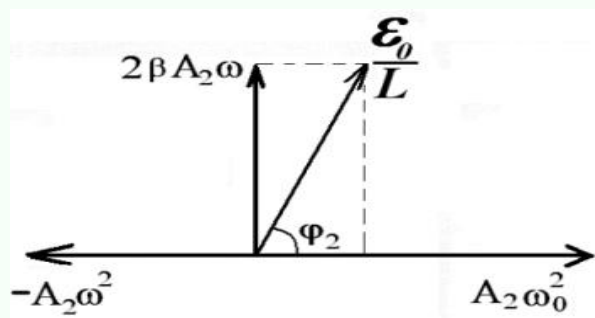


Fig.3. Vector diagram

From the vector diagram we find  $A_2$ :

$$A_2 = \frac{f_0}{\sqrt{(\omega_0^2)^2 - \omega^2 + 4\beta^2\omega^2}} = \frac{e_0/L}{\sqrt{\left(\frac{1}{LC} - \omega^2\right)^2 + 4\left(\frac{R}{2L}\right)^2 \omega^2}}$$

Here  $\omega$  is the frequency of the alternating EMF.

$$\text{tg}\varphi_2 = \frac{2\beta\omega}{\omega_0^2 - \omega^2} = \frac{R\omega}{L\left(\frac{1}{LC} - \omega^2\right)} = \frac{R}{\frac{1}{\omega C} - \omega L} \leftarrow \text{sign } \text{tg}\varphi_2 \text{ can be negative if } \omega_0 < \omega \text{ and}$$

therefore in this case the phase shift between the driving force and the displacement  $\varphi_2 > \frac{\pi}{2}$ .

$$x_2 = \frac{e_0/L}{\sqrt{\left(\frac{1}{LC} - \omega^2\right)^2 + \left(\frac{R}{L}\right)^2 \omega^2}} \cos(\omega t - \varphi_2) = q(t)$$

This solution corresponds to the law of change of charge on the capacitor plates and voltage on its plates:

$$U_C = \frac{e_0/LC}{\sqrt{\left(\frac{1}{LC} - \omega^2\right)^2 + \left(\frac{R}{L}\right)^2 \omega^2}} \cos(\omega t - \varphi_2).$$

- 1) Let us examine the obtained solution[4].
- 2) If the active resistance of the circuit  $R$  tends to 0, then the attenuation coefficient of the system:
- 3)  $\beta \rightarrow 0$  bundan  $q = A_1 \sin(\omega_0 t + \alpha_1) + \frac{e_0/L}{\frac{1}{LC} - \omega^2} \cos \omega t$

In this case, the first term does not decay,  $A_1$  and  $\alpha_1$  are determined from the initial conditions.

At  $\beta \rightarrow 0$   $A_2(\omega \rightarrow \omega_0) \rightarrow \infty$ .

- 4) Let's find the current in the circuit:

$$5) \quad I = \frac{dq}{dt} = - \frac{\frac{\omega e_0}{L}}{\sqrt{\left(\frac{1}{LC} - \omega^2\right)^2 + \left(\frac{R}{L}\right)^2 \omega^2}} \sin(\omega t - \varphi_2).$$

The current in the circuit, and therefore the voltage on the active resistance, are shifted in phase by  $\pi/2$  relative to the voltage on the capacitance and, as will be shown below, by  $(-\pi/2)$  relative to the voltage on the inductance[5].

Note that the phase shift of the current relative to the phase of the EMF is determined by the angle  $\varphi = \varphi_2 - \frac{\pi}{2}$ :

$$I = \frac{dq}{dt} = \frac{\frac{\omega e_0}{L}}{\sqrt{\left(\frac{1}{LC} - \omega^2\right)^2 + \left(\frac{R}{L}\right)^2 \omega^2}} \cos\left(\omega t - \varphi_2 + \frac{\pi}{2}\right) = \frac{\frac{\omega e_0}{L}}{\sqrt{\left(\frac{1}{LC} - \omega^2\right)^2 + \left(\frac{R}{L}\right)^2 \omega^2}} \cos(\omega t - \varphi)$$

At the same time  $\operatorname{tg}\varphi = \frac{1}{R}(\omega L - \frac{1}{\omega C})$ .

Resonance in an oscillatory circuit. Let's find the maximum of  $A_2(\omega)$  when the denominator is minimal:

$$A_2 = \frac{f_0}{\sqrt{(\omega_0^2)^2 - \omega^2 + 4\beta^2\omega^2}}$$

From this it turns out  $\omega_{\text{rez}} = \sqrt{\omega_0^2 - 2\beta^2}$  ← resonant frequency - the frequency at which an EMF with amplitude  $e_0$  can excite oscillations of voltage and charge on the capacitor plates with maximum amplitude. At  $\beta \ll \omega_0 \rightarrow R \ll \sqrt{\frac{L}{C}}$ ,

That's why  $q_{\text{max}} = q_{\text{rez}} = \frac{e_0/L}{2\beta\sqrt{-\beta^2+\omega_0^2}} \approx \frac{e_0}{R\omega_0}$ .

The obtained result can be interpreted as follows: the closer the EMF frequency is to the resonant frequency of the circuit, the greater the maximum charge that occurs on the capacitor plates[6]:

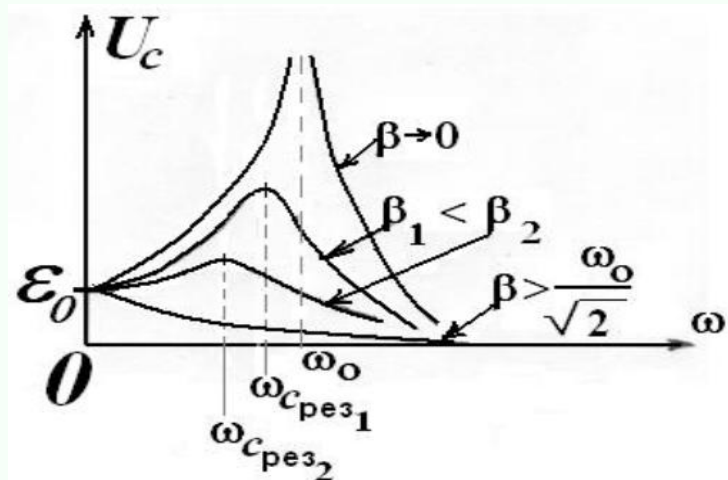


Fig.4. Graph of resonance phenomenon

From this graph it is clear that the real resonance is observed at  $\beta=0$ . When we compare electromagnetic resonance with mechanical ones, it is clear that in electromagnetic oscillatory processes there is a resonance of voltages on the capacitance, a resonance of current and also a resonance of voltages on the inductance, the conditions of which are the following[7]:

$$\omega_C(\text{rez}) = \sqrt{\omega_0^2 - 2\beta^2} = \sqrt{\frac{1}{LC} - \frac{R^2}{2L^2}}; \quad I_{\text{rez}} = \frac{e_0}{R}; \quad \omega_L(\text{rez}) = \frac{1}{\sqrt{LC - \frac{R^2C^2}{2}}}$$

The quality factor of the system, expressed through the resonance characteristics. Let the natural frequency  $\omega_0 \approx \omega_{CB}$  be the frequency of free damped oscillations. For electromagnetic oscillatory processes, the quality factor is numerically equal to the ratio  $Q = \frac{U_C(\text{rez})}{e_0}$ .

As a conclusion, it can be said that resonance in electromagnetic oscillatory circuits is used in the transmission and reception of electromagnetic waves. Since communication is mainly carried out through the ether using electromagnetic waves.

Considering that the electromagnetic spectrum includes radio waves, infrared radiation, visible radiation, ultraviolet radiation, x-rays and even gamma radiation, their military applications are as follows:

- Radio waves are used in military communications and troop control;
- Radar is used to determine the distance, speed and trajectory of movement to the target;
- Used in electronic warfare, it disrupts enemy communications, disables radars, etc.
- It is used as an optoelectronic system, such as an infrared thermal imager, a homing missile with an infrared sensor installed in the warhead of the missile, and a laser rangefinder and guidance system;
- While lasers have been used as directed energy weapons to destroy drones and as microwave radiation to disable enemy electronics, X-ray-emitting laser devices have now been developed for military use as well.

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